Phaseless Single-Step Microwave Imaging Technique for Biomedical Applications

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Abstract. In the present work, an improved phaseless approach to microwave imaging is presented. Starting from the Contrast Source formulation of the scattering problem, a single-step procedure with no intermediate phase-retrieval process is described. The reconstruction capabilities of the proposed phaseless inverse method are numerically validated by firstly considering simple dielectric targets. Then, a slice breast model with the inclusion of a cancerous portion is analyzed. The identification of different types of breast tissue is successfully achieved, thus confirming the validity and potentialities of the proposed phaseless technique in the framework of biomedical imaging.

Keywords

Electromagnetic (EM) inverse scattering problem, Phaseless Contrast-Source Inversion method (P-CSI), breast imaging

1. Introduction

Microwave imaging (MWI) techniques are gaining increasing interest in medical diagnostics, due to the adoption of non-ionizing radiation as well as the provided low Specific Absorption Rate (SAR), further combined with a cheaper, non-invasive and compact imaging setup. However, the limited penetration depth and the relatively low resolution give a strong limitation in the large-scale deployment of microwave tomography (MWT). The potentialities of MWI to provide information about the health status of inaccessible tissues, with a better dielectric representation of biological materials, suggest to combine this technique with the widely used X-ray and Magnetic Resonance Imaging (MRI). Furthermore, MWT is revealed to be less prone to false-positive and false-negative rates, as compared to conventional diagnostics techniques [1]. According to a first preliminary investigation presented in [2], a Phaseless Contrast Source Inversion Method (P-CSI) is successfully implemented in the present work for breast tissue reconstruction. In particular, a Contrast Source (CS) formulation is adopted [3], [4], and the inverse scattering

problem is solved with no linearization procedure, by converting it into an iterative optimization problem, where the two unknowns, namely the contrast source and the dielectric contrast, are alternatively updated according to a conjugate gradient scheme. The inversion procedure is performed by exploiting the amplitude-only data of the measured total field, locally defined as the sum of the incident and the scattered fields, the former obtained as a base-line measurement in the absence of the Object Under Test (OUT), and the latter due to the interaction of the incident field with the OUT. The full-data information of the incident field (which can be easily extracted from simulations) is also required for the reconstruction process.

2. Scattering Problem Formulation

Let us consider a two-dimensional tomography problem, aiming at localize a generic OUT and retrieve its dielectric properties, hereby denoted as B. A TM-polarized incident wave is considered, and cylindrical targets are analyzed as well. A magnetic permittivity equal to that of free-space is assumed, to match with the non-magnetic property of a biomedical scenario. Measurement points are equally distributed on the acquisition curve S, as shown in the general scheme of Fig. 1.

It describes a multi-static and multi-view setup, where the transmitter location is alternatively changed, with N_{TX} angles of incidence, resulting in a number $N_{\text{RX}} \times N_{\text{TX}}$ of



Fig. 1. Configuration of acquisition setup.

measurements. The imaging domain, also named Domain of Interest (DOI) and hereby indicated as ${\rm I\!D}$, fully contains the unknown OUT.

According to the theory [5], [6], the scattering problem can be analyzed by considering the following equations, known as Electrical Field Integral Equations (EFIEs):

$$E^{t}(r) = E^{i}(r) + k_{b}^{2} \int_{\mathbb{D}} G(r, r') \chi(r') E^{t}(r') dr', r \in \mathbb{S}, (1)$$

$$E^{\mathrm{s}}(r) = k_{\mathrm{b}}^{2} \int_{\mathbb{D}} G(r, r') \chi(r') E^{t}(r') \mathrm{d}r', r \in \mathbb{D}.$$
 (2)

In particular, equation (1), indicated as *data equation*, relates the measured scattered field along the acquisition domain S with the dielectric contrast, conventionally expressed in a normalization form with respect to the background permittivity and indicated as χ :

$$\chi(r) = \frac{k^2}{k_{\rm b}^2} - 1 \tag{3}$$

 $k_{\rm b}$ indicating the background wavenumber.

Equation (2), known as *state equation*, relates the total field, evaluated inside the DOI, with the dielectric properties inside the domain \mathbb{D} , wherein the object \mathbb{B} is assumed to be delimited.

Let us assume V different positions given by a TM source; for a fixed position, hereby indicated by v, the measured field is sampled along the acquisition curve by changing the location of the receiving antenna. Thus, series of measurement campaigns are performed. According to the aforementioned notation, the EFIEs can be reformulated in a compact form, by introducing the contrast source term, namely $\omega(r) = \chi(r)E^{t}(r)$. We can write then:

$$E_{\nu}^{t}(r) = E^{i}(r) + G_{s}\omega_{\nu}(r), \quad r \in \mathbb{S}, \qquad (4)$$

$$\omega_{v}(r) = \chi(r)[E_{v}^{i}(r) + G_{D}\omega_{v}(r)], \ r \in \mathbb{D}$$
 (5)

where G denotes the same radiation operator for both equations, with the respective limitation in the range of r and easily expressed in terms of Hankel function of the zero-th order and second kind.

In the Contrast Source Imaging (CSI) method, the solution of the tomographic problem requires the definition of a quadratic cost functional to be minimized. It is obtained as the combination of two sub-functionals, related to the misfit between measured and computed data, according to (4), (5).

From the above considerations, it results:

$$F(\omega_{v}, \chi) := F_{S}(\omega_{v}) + F_{D}(\omega_{v}, \chi)$$
(6)

where F_S and F_D are denoted as *data function* and *state function*, respectively, and they can be expressed as:

$$F_{\rm S}(\omega_{\rm v}) = \alpha_{\rm S} \sum_{\nu} \left\| \left| f_{\nu} \right|^2 - \left| E_{\nu}^{\rm i} + G_{\rm S} \omega_{\nu} \right|^2 \right\|_{\rm S}^2, \qquad (7)$$

$$F_{\rm D}(\omega_{\nu}) = \alpha_{\rm D} \sum_{\nu} \left\| \chi(E_{\nu}^{\rm i} + G_{\rm D}\omega_{\nu}) - \omega_{\nu} \right\|_{\mathbb{D}}^{2}$$
(8)

where α_{S-D} are proper normalization factors, while f_v indicates the measured total field along the acquisition curve.

As it can be easily observed from (7), (8), amplitudeonly data for the total field are exploited, while the incident field is assumed in both amplitude and phase. A suitable incident field inside the DOI is obtained by looking for a computed matching incident field along the acquisition curve first [7], [8]. The latter is obtained as a linear combination of a limited number N of Hankel functions of the second kind, so that the following relation holds true:

$$E^{i}_{s}(\rho,\theta) \cong -\frac{j}{4} \sum_{\nu=-N}^{N} c_{\nu} H_{\nu}^{(2)}(k\rho) \exp(j\nu\theta) \qquad (9)$$

where (ρ, ϑ) stands for polar coordinates.

In the second step, coefficients c_v are re-used for the determination of the incident field inside \mathbb{D} [9].

The reconstruction procedure requires the CS initialization. The choice for this initial value into the optimization algorithm results to be a key point for the imaging process, since the convergence of the local optimization problem can be strongly influenced by the above initial guess.

According to the original work [3] introducing the CSI method, the use of the back-propagation (BP) approach for the initial estimate for the CS requires the fulldata availability of the total field. Consequently, further techniques have been proposed in literature with the aim to exploit the BP method in a phaseless approach, by considering an arbitrary phase for the total field [10], [11]. However, due to the lack of reason behind the phase selection, an alternative solution for the CS initialization is considered in the present work. This is obtained by applying the steepest descent method to the data function F_S , limited to the first step. In accordance with the aforementioned approach, it consequently results:

$$\omega_{\nu,0} = -2\beta\alpha_{\rm S} G^*_{\rm S} \left[E^{\rm i}_{\nu} \left(\left| f_{\nu} \right|^2 - \left| E^{\rm i}_{\nu} \right|^2 \right) \right]$$
(10)

where the symbol * denotes the adjoint operator, while parameter β indicates the step size. According to (10), all available data can be fully exploited into the initialization process; furthermore, no a priori information about the contrast inside \mathbb{D} is required, as compared to the approach outlined in [12], where an approximate knowledge of the average contrast over the DOI is necessary.

3. Numerical Validations on Dielectric Targets Reconstruction

A first numerical validation of the proposed method is performed by using available measured data from Fresnel



Fig. 2. Permittivity map for the target FoamDielInt in [13].

FoamDielInt	٤ _r	Min	Mean	Max
Air	1	0.97	1.00	1.08
Foam	1.45±0.15	0.98	1.29	1.86
Plastic	3±0.3	2.27	2.43	2.78

Tab. 1. Reconstructed permittivity values (a) f= 2 GHz, with a maximum number of iterations equal to 150.

Institute. They are relative to cylindrical dielectric targets combined in different fashions, thus resulting in homogeneous as well as inhomogeneous targets [13], [14].

The algorithm is tested for a fixed operating frequency, while the reconstruction process is performed until a convergence criterion on the cost functional $F(\omega_v, \chi)$ is met or, alternatively, when a fixed maximum number of iterations is exceeded. An example of dielectric reconstruction is shown in Fig. 2, whereas the retrieved permittivity values are indicated in Tab. 1.

4. Breast Tissue Modeling and Imaging Results

In order to confirm the reconstruction capabilities of the proposed method for breast imaging applications, a slice breast model is implemented on COMSOL Multiphysics[®] platform [15]. A simplified 3-tissue breast model is considered, in which a malignant portion with a 5 mm radius is included into the adipose tissue; the respective dielectric properties are determined similarly to [16] for a 2 GHz operating frequency and are listed in Tab. 2, while a matching medium with $\varepsilon_{rb} = 12$ is assumed.

The detection of the tumor location comes from the permittivity contrast due to the different water content of the malignant tissues compared to the fat, combined with

Tissue	٤ _r	σ
Fat	5	0.05
Fibroglandular tissue	43	0.7
Tumor	35	0.7

Tab. 2. EM properties of breast tissues.



Fig. 3. Slice breast model on COMSOL platform.



Fig. 4. Computational mesh and PML on COMSOL platform.

the higher conductivity of the cancerous portions. The breast for the 2D modeling is emulated with a cylinder having a 45.5 mm radius, while the fibroglandular portion has a radius equal to 20 mm. The computational domain filled by the background medium needs to be truncated with the definition of a perfectly matched layer (PML) as boundary condition (Fig. 3). The correspondent computational mesh is shown in Fig. 4. Point sources are implemented for the TM field generation, and thus a series of line current out-of-plane are considered on COMSOL platform, in order to speed-up the forward computation needed for the generation of the synthetic data.

As shown in Fig. 5, a quantitative reconstruction of the contrast function χ and, consequently, of the relative permittivity is obtained for the lower-contrast regions. In particular, the cancerous region is clearly localized in the right corner.

The current state of the algorithm allows to retrieve a full-dielectric map in the case of relatively low-contrast scenario. In the case of high permittivity contrast, a qualitative reconstruction is guaranteed; additional developments are currently performed to optimize the performances.

5. Conclusion

The potentialities of a phaseless microwave tomography approach for biomedical imaging applications has been described and discussed in this work. Firstly, the numerical



Fig. 5. Real (a) and imaginary (b) part of the dielectric contrast χ for the slice breast model.

validation of the proposed method has been performed by the reconstruction of a series of dielectric target. Secondly, a permittivity map of a slice breast has been obtained, able to localize a malignant tissue inside the adipose region of the breast model. The discussed results represent a useful preliminary assessment for the implementation of a lowcost, compact and non-ionizing breast imaging setup, able to support the currently adopted diagnostic techniques in the framework of breast cancer detection and breast imaging applications.

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