

# Miniaturization of Power Dividers by Using Asymmetric CMRC Structures and QWLTs with Low-Cost Materials

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**Abstract.** *This paper presents the miniaturization of power dividers using asymmetric compact-microstrip-resonant-cell (CMRC) structures employing low-cost materials based on a quarter-wave-like transformer (QWLT). The proposed CMRC-based QWLT power divider is intended for operation at a frequency of 2.4 GHz, utilizing the FR-4 print circuit board (PCB) with a dielectric constant of 4.3 and a substrate thickness of 1.6 mm. The CMRC dimensions include a width of 5.32 mm and a length of 8.52 mm. It is found that a significant 50% size reduction of length is achieved compared to a conventional power divider, while maintaining an insertion loss (IL) of 3.3 dB, as well as achieving the return loss and isolation loss of 20 dB.*

## Keywords

Power divider, quarter-wave transformer (QWT), quarter-wave-like transformer (QWLT), compact microstrip resonant cell (CMRC)

## 1. Introduction

The size of a conventional power divider (PD) depends on an operating frequency [1]. Thus, the traditional PD configuration is cumbersome and unsuited for modern communication systems. A multi-frequency PD has been proposed to enhance the performance of PDs, such as increasing bandwidth and miniaturizing their configuration. In this design, stubs are integrated to broaden the frequency band [2]. In addition, the bandwidth control is achieved through the utilization of diode technology and a DC power source [3]. Furthermore, a coupled line configuration in [4] is incorporated to further enhance the bandwidth. The techniques presented in [2–7] were employed to decrease losses of PDs. However, the physical dimensions of PDs in those prior studies remained unchanged. Various research works have been proposed to address the challenge of reducing the size of PDs.

In [8], an effective miniaturization technique of quarter-wave transformers (QWTs) was proposed based on the quarter-wave-like transformer (QWLT) theory. Note that QWLTs are implemented using multi-section transmission lines (TLs). In addition, multi-section QWLT characteristics are derived analytically and solved via appropriate optimization algorithms for associated TL parameters. It is found that QWLT prototypes yield acceptable return loss without significant bandwidth reduction, comparing to the QWT result.

Previous research works have focused on minimizing the PD size through the utilization of cost-effective FR-4 print circuit boards (PCBs) [9], [10]. In [9], the approach involves the application of multi-section transmission lines for implementing QWLTs in PDs, leading to a significant 25% size reduction compared to that of the traditional PD configuration. An alternative method for size reduction involves the implementation of a compact microstrip resonant cell (CMRC) [11], yielding a total size reduction of 40% in comparison to that of the conventional TL structure [11–14].

Several types of symmetric CMRCs that differing in shape, dimension, and transmission characteristics were proposed. In [15], this proposed configuration facilitates a 33.6% area reduction of PDs. On the other hand, as presented in [16], the miniaturization of PDs exceeds 50% in terms of area. It is noted that, similar to [15], this approach also involves three-element chips. However, installing and soldering the chip elements onto PDs are challenging due to the tiny area allocated for the installed chip elements. In the approach presented in [17], employing an open stub resulted in a size reduction of over 30% for PDs at the frequency of 1.65 GHz. This solution employs a single chip element and a simple structure for fabrication. As reported in [18], the finite defected ground structure (FDGS) achieves a size reduction of 32% for PDs at the frequency of 2.5 GHz. The solution outlined in [19] employs the substrate integrated waveguide (SIW) to decrease the size of PDs, resulting in a size reduction of 10%.

The studies conducted from [16] to [18] demonstrate in the miniaturization of PDs. These solutions exhibit the potential to reduce the PD area ranging from 30% to 52%. In the case of [16], it employs a simple structure with a cost-effective material, incorporating two chip elements in the circuit. Moreover, this solution achieved a smaller size reduction of 33.61% compared to other approaches. The research works in [17] and [18] achieved the PD miniaturization of 32–35% using a single chip element. It is noted that the references [17–19] utilize RT/Duroid and Roger materials, which are relatively expensive. Some previous research works are complicated in the fabrication and installation of chip elements and measurements. In contrast, in terms of a size reduction of conventional PDs, these previous studies achieved a size reduction lower than 40% at the operating frequency. For these reasons, this paper proposes the miniaturization of PDs by employing the QWLT-based PD introduced in [9]. The proposed approach is achieved by modifying the structure from a symmetric CMRC [11] to an asymmetric CMRC configuration in order to further decrease the size of associated TLs for PDs.

The rest of this paper is organized as follows. Section 2 presents brief theories of QWLTs and PDs. In addition, Section 3 shows results and discussions of the proposed CMRC-based QWLT PDs. Finally, conclusions are provided in Sec. 4.

## 2. Brief Theories of QWLTs and PDs

In this section, the theories of QWLTs and PDs are presented in brief, where the QWLT theory is given first.

### 2.1 Brief QWLT Theory

In [10], the derivation and properties of QWLTs are presented in details. The theory of QWLTs is briefly summarized below. Figure 1(a) illustrates the QWLT matching circuit, where  $Z_{in}$  and  $Z_L$  denote the input and load impedances respectively, where  $Z_{in}$  and  $Z_L$  are real. In addition, Figure 1(b) shows the QWLT model, where  $Z_0^+$  and  $Z_0^-$  are the characteristic impedances of the QWLT for the forward and backward propagation directions respectively, and  $\theta_c$  is the electrical length of the QWLT. Note that  $Z_0^\pm$  are complex conjugates of each other, which can be readily derived from the associated  $ABCD$  parameters of actual circuits to implement QWLTs [10], which can be written compactly as

$$Z_0^\pm = \frac{\mp 2B}{A - D \mp j\sqrt{4 - (A + D)^2}} \quad (1)$$

where  $Z_0^\pm$  can be readily expressed in the polar form as

$$Z_0^\pm = |Z_0^\pm| e^{\mp j\phi}. \quad (2)$$

Note that  $|Z_0^\pm|$  and  $\phi$  in (2) are the magnitude and argument of  $Z_0^\pm$  respectively, and are defined in the following range:

$-90^\circ \leq \phi \leq 90^\circ$ . In addition, the effective propagation constant  $\beta$  of the QWLT can be obtained from its  $ABCD$  parameters as [10]

$$\cos \theta_c = \frac{A + D}{2} \quad (3)$$

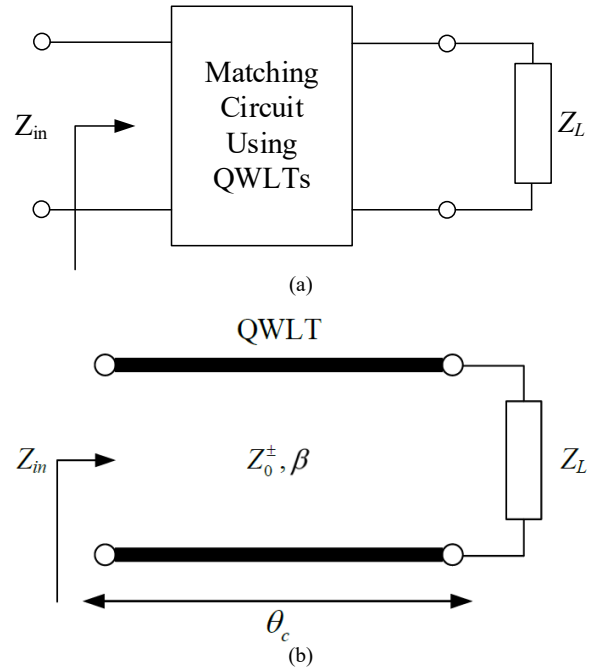
where  $\theta_c = \beta d$  and  $d$  is the physical length of the QWLT. By definition, the slow-wave factor ( $SWF$ ) of the QWLT depends on the wavelength  $\lambda_0$  of a conventional TL and the wavelength  $\lambda$  of a slow-wave TL structure as follows:

$$SWF = \frac{\lambda}{\lambda_0} \quad (4)$$

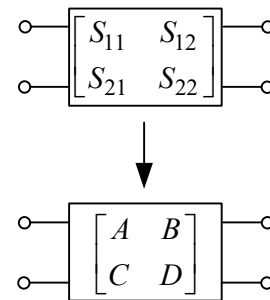
where  $\lambda$  can be determined from  $\beta$  defined in (3) as

$$\lambda = \frac{2\pi}{\beta}. \quad (5)$$

Note that a higher  $SWF$  results in a more size reduction of TL structures compared to that of the conventional TL.



**Fig. 1.** (a) The QWLT matches an arbitrary real load impedance  $Z_L$  to a specific real input impedance  $Z_{in}$ . (b) The QWLT model.



**Fig. 2.** Conversion of the scattering parameters to the  $ABCD$ -parameters for the QWLT-parameter calculation.

In [14], it is shown that the electrical length  $\theta_c$  of the QWLT is given by

$$\theta_c = \begin{cases} \tan^{-1}(-\Gamma_{f,L} \cot \phi), & \text{when } Z_L > Z_{in} \text{ and } \phi < 0, \text{ or} \\ & Z_L < Z_{in} \text{ and } \phi > 0 \\ \tan^{-1}(-\Gamma_{f,L} \cot \phi) + \pi, & \text{when } Z_L > Z_{in} \text{ and } \phi > 0, \text{ or} \\ & Z_L < Z_{in} \text{ and } \phi < 0 \end{cases} \quad (6)$$

where the reflection coefficient  $\Gamma_{f,L}$  in (6) is defined as

$$\Gamma_{f,L} = \frac{Z_L - Z_{in}}{Z_L + Z_{in}}. \quad (7)$$

In (7),  $Z_{in}$  and  $Z_L$  are both real values and given for a specific impedance-matching problem. In addition, it is shown in [14] that  $|Z_0^\pm|$  can be compactly expressed as

$$|Z_0^\pm| = \sqrt{Z_{in} Z_L}. \quad (8)$$

In this paper, the asymmetric CMRCs are implemented by using asymmetric structures with the required conditions of QWLTs given in (6) to (8), where the  $ABCD$ -parameters in computing  $Z_0^\pm$  and  $\beta$  of QWLTs can be obtained from the associated scattering parameters of QWLTs using the standard conversion of two-port parameters as shown in Fig. 2 [1].

To design a QWLT for given  $Z_{in}$  and  $Z_L$ , it is more convenient to employ the following formulation to obtain valid  $ABCD$  parameters of actual circuits to implement the desired QWLT, and then compute the QWLT parameters from the known  $ABCD$  parameters. Note that the relationship between  $Z_{in}$  and  $Z_L$  in Fig. 1(a) can be written as follows [8]:

$$Z_{in} = \frac{AZ_L + B}{CZ_L + D} \quad (9)$$

where  $A$ ,  $B$ ,  $C$ , and  $D$  in (9) are the  $ABCD$ -parameters of the QWLT matching circuit. Using (9) under the condition of reciprocal QWLTs ( $AD - BC = 1$ ) and real  $Z_{in}$  results in the following simple relationships [8]:

$$\frac{A}{D} = \frac{Z_{in}}{Z_L}, \quad (10)$$

$$\frac{B}{C} = Z_{in} Z_L. \quad (11)$$

Note that the required conditions in (10) and (11) are more convenient to solve and easy to calculate associated QWLT parameters via (1) and (3).

## 2.2 Brief Theory of PDs

In this section, the theory of both conventional and CMRC-based QWLT PDs is discussed in brief.

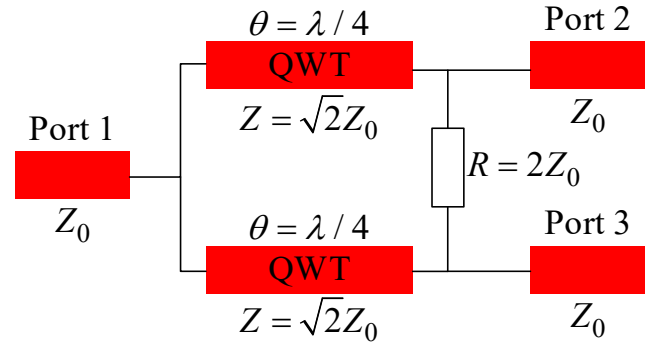


Fig. 3. The schematic diagram of the conventional PD.

### 2.2.1 Conventional PDs

The conventional PD is designed by using two standard QWLTs [1]. Figure 3 shows the schematic diagram of the conventional PD, where  $\lambda$  is the operating wavelength,

$$\theta = \lambda/4, \quad (12)$$

$$Z = \sqrt{2}Z_0, \quad (13)$$

$$R = 2Z_0. \quad (14)$$

Note that  $\theta$  is the electrical length of QWLTs,  $Z$  is the characteristic impedance of QWLTs,  $R$  is the isolation impedance of the conventional PD, and  $Z_0$  is the port impedance. Conventional PDs are simple to design and fabricate; however, their dimensions are determined by the operating wavelength, usually resulting in larger sizes for lower frequencies. In addition, traditional PDs cannot be readily changed in their shape or reduced in size.

### 2.2.2 CMRC-based QWLT PDs

This subsection shows the concept of CMRC-based QWLT PDs. Figure 4 illustrates the schematic diagram of CMRC-based QWLT PDs, where  $Z_0^\pm$  are the characteristic impedances of QWLTs,  $\theta_c$  is their electrical length and  $Z_{S,CMRC}$  is the isolation impedance. Note that  $Z_0^\pm$  and  $\theta_c$  can be readily determined using the even-odd mode analysis [1], [9] and  $Z_{S,CMRC}$  is readily calculated based on the principle outlined in [9]. It should be noted that, with this approach, the CMRC-based QWLT for the even-mode analysis provides that the load impedance  $Z_L = 2Z_0$  is matched to the input impedance  $Z_{in} = Z_L$ . Based on these impedances and a certain value of the argument  $\phi$  defined in (2), the results of minimizing the divider design depending on the QWLT parameters are calculated according to (6) to (8). As pointed out in [9], the  $ABCD$ -parameters of the QWLT in terms of  $Z_0^\pm$  and  $\theta_c$  are useful in the derivation of  $Z_{S,CMRC}$  as given below:

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix} = K \begin{bmatrix} Z_0^+ e^{j\theta_c} + Z_0^- e^{-j\theta_c} & Z_0^+ Z_0^- (e^{j\theta_c} - e^{-j\theta_c}) \\ e^{j\theta_c} - e^{-j\theta_c} & Z_0^- e^{j\theta_c} + Z_0^+ e^{-j\theta_c} \end{bmatrix}, \quad (15)$$

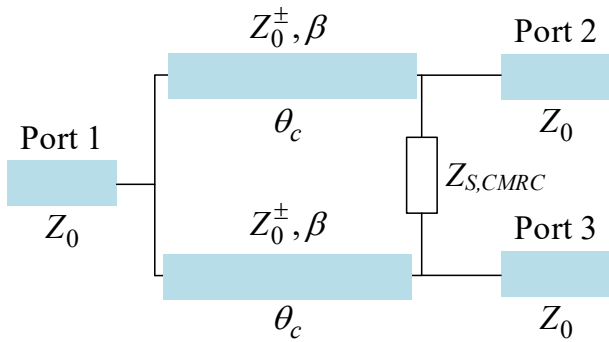


Fig. 4. The schematic diagram of the CMRC-based QWLT PDs.

$$K = \frac{1}{Z_0^+ + Z_0^-}. \tag{16}$$

In the next section, associated results of both conventional and CMRC-based QWLT PDs are provided. In addition, comparisons and discussions are also given.

### 3. Results and Discussions

The core concept of this paper involves modifying a symmetric CMRC to an asymmetric CMRC through an application of QWLTs. Figures 5 and 6 illustrate the symmetric and asymmetric CMRC structures, respectively. The goal of this study is to achieve a 50% size reduction at the operating frequency of 2.4 GHz. Note that the prototype of the proposed CMRC-based QWLT PD is shown in Fig. 7(b), where its structure is compact and simple for implementation. The conventional and proposed PDs are designed on a lossless FR-4 substrate material with a dielectric constant of 4.2 and a thickness of 1.6 mm. The conductor is assumed to be a perfect electric conductor (PEC) to ensure lossless TL characteristics. This study compares the performance of the two PDs at 2.4 GHz. Note that the isolation resistance  $R_S$  in Fig. 7(a) for the conventional Wilkinson PD is obtained as shown in [1], while the isolation impedance  $Z_{S,CMRC}$  in Fig. 7(b) for the proposed CMRC-based QWLT PD is designed according to [9].

The conventional Wilkinson PD is shown in Fig. 7(a), where the length of each TL is equal to  $\lambda/4$  (QWT) and the



Fig. 5. A symmetric CMRC structure [11].

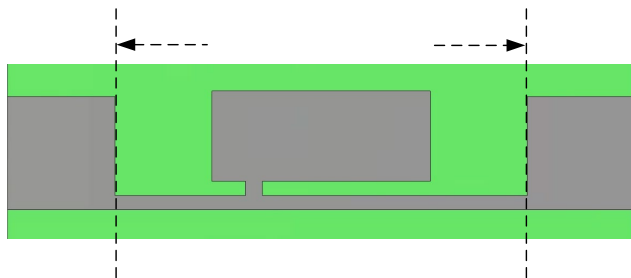
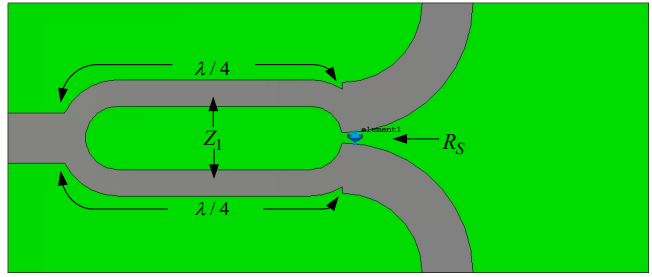
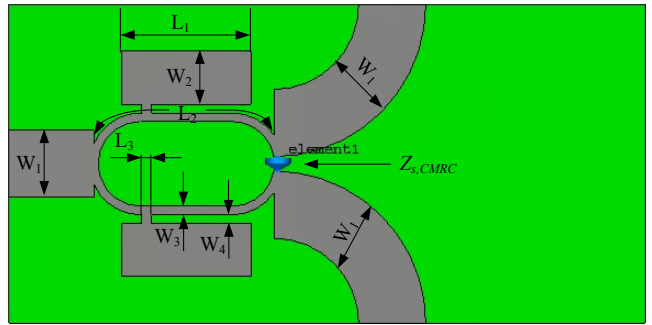


Fig. 6. An asymmetric CMRC structure modified from the structure in Fig. 5.



(a)



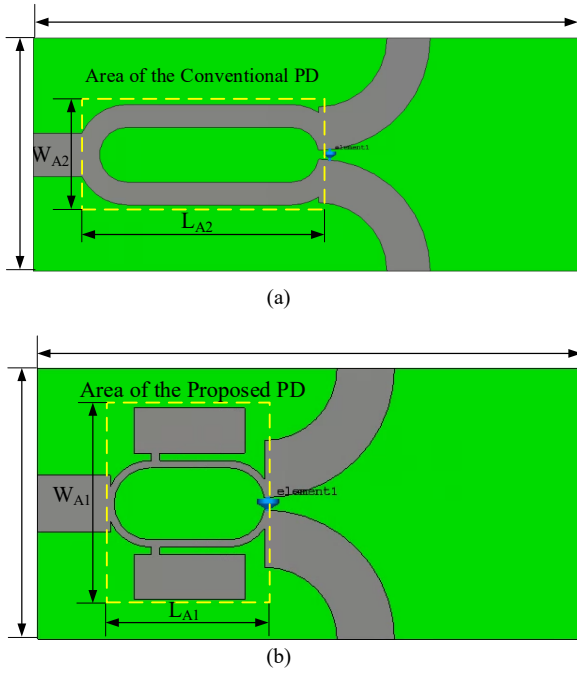
(b)

Fig. 7. (a) Conventional Wilkinson PD. (b) The proposed CMRC-based QWLT PD.

Parameter	Dimension (mm)
$W_1$	3.16
$W_2$	2.51
$W_3$	0.40
$W_4$	0.40
$L_1$	6.08
$L_2$	11.48
$L_3$	0.48

Tab. 1. Dimensions of the CMRC-based QWLT PD.

isolation impedance  $R_S$  is given in (14). In addition, the characteristic impedance  $Z_1$  of two QWTs is provided in (13). Note that  $\lambda/4$  is equal to 22.96 mm at 2.4 GHz in Fig. 7(a) and the length  $L_2$  of the proposed CMRC-based QWLT PD is equal to 11.48 mm in Fig. 7(b). The QWLT based on an asymmetric CMRC structure in Fig. 7(b) can be found using the even-mode analysis as detailed in [1]; i.e., the QWLT transforms the load impedance of  $2Z_0$  to the input impedance of  $Z_0$ , where  $Z_{in}$  and  $Z_L$  in (10) and (11) in this case are equal to  $Z_0$  and  $2Z_0$ , respectively. Note that the  $ABCD$  parameters in (10) and (11) are determined numerically by simulation of the proposed asymmetric CMRC structure in Fig. 6 using the CST-MW Studio [20] for desired characteristics. By solving (10) and (11) simultaneously to obtain the  $ABCD$  parameters, each optimized asymmetric CMRC structure is shown in Fig. 7(b), where its parameters are given in Tab. 1. Once the  $ABCD$  parameters are known, the QWLT parameters can be computed using (1) to (3) and each QWT of the conventional Wilkinson PD can be replaced by the QWLT as shown in Fig. 7(b) [11].



**Fig. 8.** The circuit area and dimension of (a) the conventional Wilkinson PD and (b) the proposed CMRC-based QWLT PD.

In [9], after straightforward manipulation,  $Z_{S,CMRC}$  can be expressed compactly as

$$Z_{S,CMRC} = 2Z_0 \frac{\sin^2 \theta_c}{\cos^2 \phi} - j\sqrt{2}Z_0 \frac{\sin 2\theta_c}{\cos \phi} \quad (17)$$

where  $\theta_c$  of each QWLT is equal to  $69.26^\circ$  (using (6)),  $\phi$  of each QWLT in Fig. 7(b) and using Tab. 1 is chosen to be  $7.2^\circ$  [9], and  $|Z_0^\pm| = 70.71 \Omega$  (using (8) with  $Z_{in} = Z_0$  and  $Z_L = 2Z_0$ ). Using (17), it is found that  $Z_{S,CMRC}$  is equal to  $(88.86 - j47.20) \Omega$ , which can be implemented using the resistor of  $88.70 \Omega$  and the capacitor of  $1.45 \text{ pF}$  at  $2.4 \text{ GHz}$ .

Figures 8(a) and 8(b) show the circuit dimension of the conventional Wilkinson PD and the proposed CMRC-based QWLT PD. In addition, Table 2 presents a comparison of the circuit dimension between two PDs. It is found that the total area of the conventional Wilkinson PD in Fig. 8(a) is equal to  $679.20 \text{ mm}^2$ , while that of the proposed CMRC-based QWLT PD in Fig. 8(b) is equal to  $450.30 \text{ mm}^2$ . In addition, the PD areas of the conventional Wilkinson PD and proposed CMRC-based QWLT PD as shown in Fig. 8 are equal to  $126.87 \text{ mm}^2$  and  $90.35 \text{ mm}^2$ , respectively. Thus, it is obvious that both areas of the proposed PD are significantly smaller than those of the conventional one.

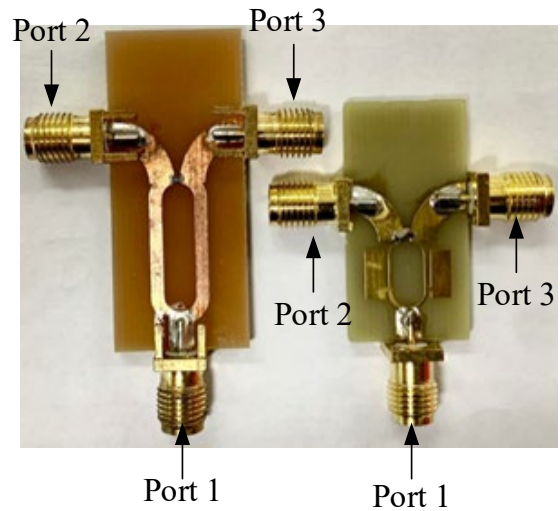
Table 3 shows the QWLT parameters of the asymmetric CMRC structure in Fig. 7(b), computed using (1) to (8) by the  $ABCD$ -parameters obtained as a result of converting the optimized scattering parameters. The magnitudes of characteristic impedances  $|Z_0^+|$  and  $|Z_0^-|$  are equal to  $77.62 \Omega$  and  $77.64 \Omega$  respectively, which are slightly different. In addition, their phases,  $\phi^+$  and  $\phi^-$ , are equal to  $-8.65^\circ$  and  $8.60^\circ$ , respectively. Note that  $Z_0^+$  and  $Z_0^-$  are almost complex conjugate of each other as expected.

Parameter	Dimension (mm)
$L_{S1}$	30
$L_{S2}$	40
$L_{A1}$	8.52
$L_{A2}$	17.38
$W_{S1}$	15.01
$W_{S2}$	16.98
$W_{A1}$	10.63
$W_{A2}$	7.30

**Tab. 2.** Dimension comparison between the conventional Wilkinson PD and the proposed CMRC-based QWLT PD.

QWLT parameter	Value
$ Z_0^+ $	$77.62 \Omega$
$ Z_0^- $	$77.64 \Omega$
$\phi^+$	$-8.65^\circ$
$\phi^-$	$8.60^\circ$
SWF	3.66

**Tab. 3.** The QWLT parameters of the asymmetric CMRC structure in Fig. 7(b) using Tab. 1.



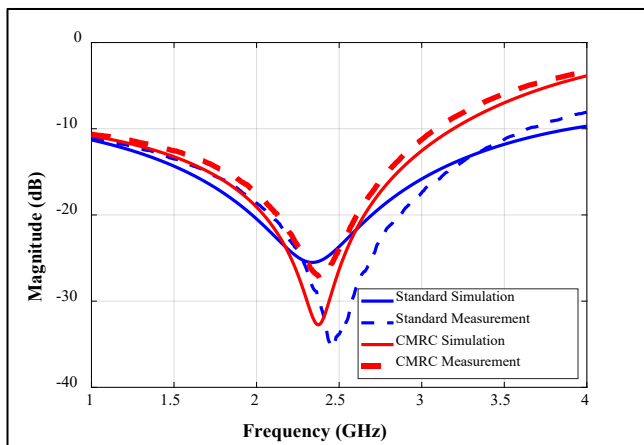
**Fig. 9.** The fabricated prototypes of the conventional Wilkinson PD (left) and the proposed CMRC-based QWLT PD (right).

Figure 9 shows the fabricated prototypes of the conventional Wilkinson PD and the proposed CMRC-based QWLT PD. From Fig. 9, the size comparison reveals that the PD length of the latter is about 50% of that of the former, where the latter provides a significant size reduction indeed.

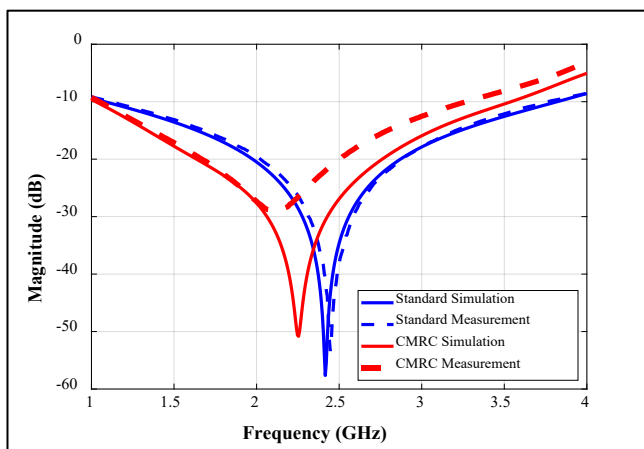
Figure 10 illustrates the simulated and measured frequency responses of the CMRC-based QWLT PD compared to those of the conventional Wilkinson PD. Note that Port 1 is the input port, while Ports 2 and 3 represent the output

ports. Note that Figure 10(a) shows the magnitude of  $S_{11}$  ( $|S_{11}|$ ), while Figure 10(b) illustrates the isolation results ( $|S_{23}|$ ). It is found that  $|S_{11}|$  and  $|S_{23}|$  are below  $-25$  dB and  $-20$  dB at 2.4 GHz, respectively. It is found that the simulated and measured results are in good agreement. However, there is a slight discrepancy in resonant frequencies between the measured and simulated results for both conventional and CMRC-based QWLT PDs. In addition, Figure 11 shows a comparison of  $|S_{21}|$  and  $|S_{31}|$  of simulated and measured results for both conventional and CMRC-based QWLT PDs. It is found that the slight discrepancy in magnitude of about 0.5 dB at 2.4 GHz between simulated and measured results is obtained.

In addition, Table 4 compares the proposed CMRC-based QWLT PD with other PDs selected from previous research works, emphasizing key parameters such as the length reduction, the number of chip elements used, the substrate material, the operating frequency, the circuit size, and the S-parameters given at the lowest frequency for each reference. Note that the PDs in [2–7] did not achieve any size reduction. Furthermore, PDs in [3], [6], and [7] utilized both resistor and capacitor chips, leading to complex structures. In [9] and [15–19], the size reduction of PDs can be achieved as shown in Tab. 4. Note that the PD in [16] incorporated

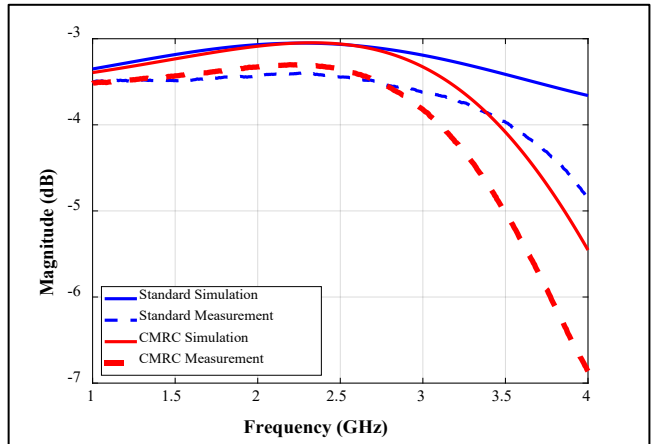


(a)

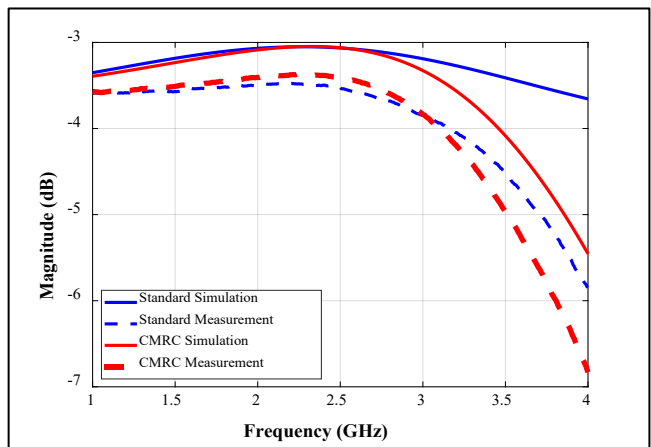


(b)

**Fig. 10.** Simulated and measured frequency responses of the conventional Wilkinson PD and the CMRC-based QWLT PD: (a)  $|S_{11}|$ , (b)  $|S_{23}|$ .



(a)



(b)

**Fig. 11.** Simulated and measured insertion losses of the conventional Wilkinson PD and the CMRC-based QWLT PD: (a)  $|S_{21}|$ , (b)  $|S_{31}|$ .

more than two chip elements in the circuit, and the PD in [19] has no chip element but its size reduction is lower than 20%.

To demonstrate the flexibility in controlling the passband by adjusting the center frequency, the total length of the CMRC-based QWLT PD ( $L_{A1}$ ) and the total length of the conventional Wilkinson PD ( $L_{A2}$ ) defined in Fig. 8 are varied to see their effects on the frequency responses of  $|S_{11}|$ ,  $|S_{21}|$ ,  $|S_{31}|$  and  $|S_{23}|$  as shown in Fig. 12. In this paper, both PDs are designed at 2.4 GHz, where  $L_{A1}$  and  $L_{A2}$  are equal to 8.52 and 17.38 mm (the designed values) as shown in Tab. 2, respectively. It is found that as  $L_{A1}$  and  $L_{A2}$  are decreased from the designed values, the center frequency is shifted towards the higher frequency and vice versa as expected. Thus, the CMRC-based QWLT PD is flexible to control the passband by adjusting the center frequency indeed. Finally, it should be pointed out that the described approach in this paper can be applied to systematically design similar power dividers and other microwave circuits if they possess the quarter-wave transformers (QWTs) in their circuits by replacing the QWTs by the corresponding QWLTs based on asymmetric CMRC structures. Note that the QWLTs can be systematically designed by using the brief design procedure outlined in Sec. 3 (see the description of Fig. 7(b)).

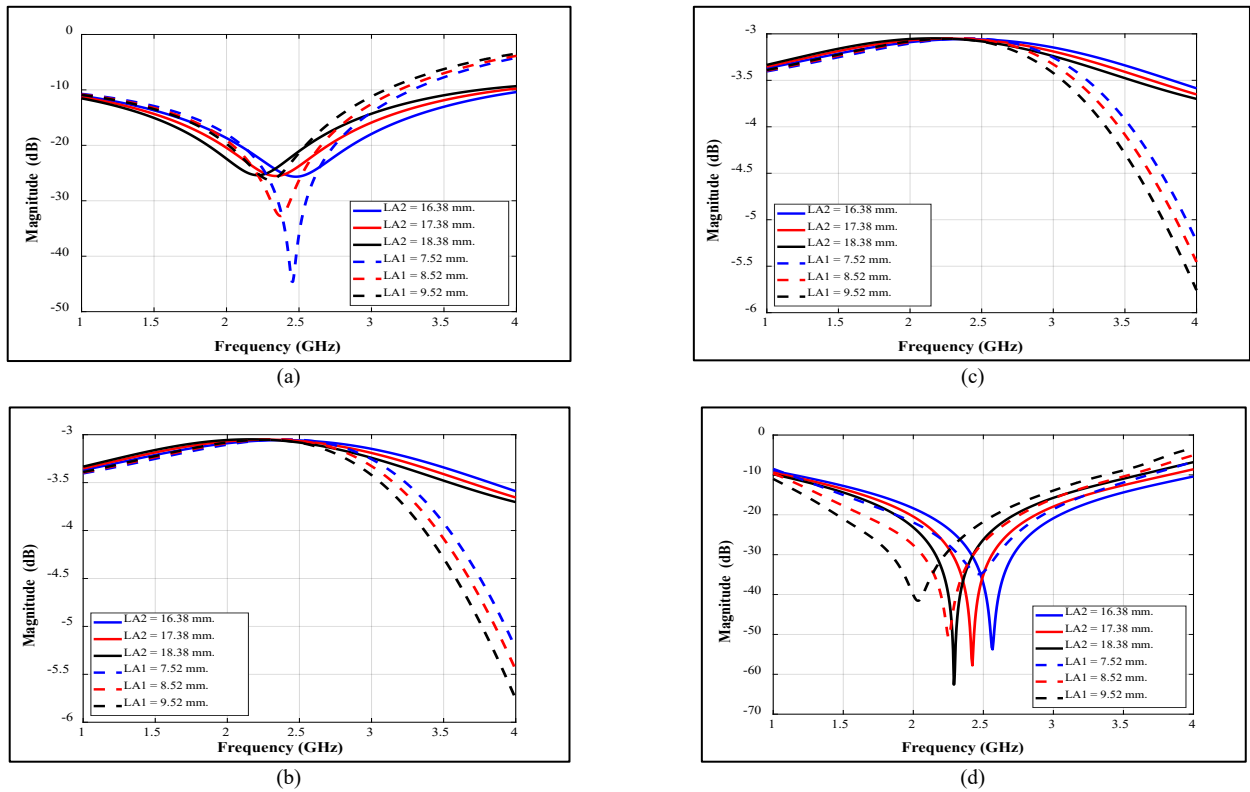


Fig. 12. Effects of  $L_{A1}$  and  $L_{A2}$  to the frequency responses of the CMRC-based QWLT PD and the conventional Wilkinson PD, respectively: (a)  $|S_{11}|$ , (b)  $|S_{21}|$ , (c)  $|S_{31}|$ , (d)  $|S_{23}|$ .

Ref.	Substrate material	Number of chip elements	Length reduction compared with $\lambda_g/4$ (%)	Operating frequency (GHz)	Circuit size ( $\lambda_g \times \lambda_g$ )	$ S_{11} ,  S_{21} ,  S_{31} ,  S_{23} $ (dB)
[2]	RO4350B	0	0	2.0, 4.4 and 5.0	$2.15 \times 2.15$	-20.00, -4.44, -3.10, -7.39
[3]	RO4350B	11	0	2.4	$2.03 \times 1.37$	-20.00, -1.12, -14.90, -13.70
[4]	RT/Duroid 5880	1	0	1.0-7.0	$1.96 \times 2.09$	-36.00, -3.24, -3.24, -29.00
[5]	RT/Duroid 5880	0	0	14.5	$2.55 \times 1.68$	-33.00, -3.20, -3.20, -22.00
[6]	FR-4	9	0	3.5	$3.41 \times 2.56$	-25.00, -3.00, -3.00, -52.00
[7]	Material (dielectric constant of 2.94)	1,5 and 8	0	10.5	$2.11 \times 2.11$	-21.00, -3.30, -3.30, -20.00
[9]	FR-4	2	25	2.4	$0.77 \times 0.74$	-22.00, -3.60, -3.60, -30.00
[15]	FR-4	2	0	2.4	$1.74 \times 1.51$	-26.73, -3.67, none, none
[16]	Roger 4003C	3	0	2.4, 5.9	$0.73 \times 0.73$	-26.10, -3.05, -3.30, -28.00
[17]	RT/Duroid 5880	1	28.76	1.65	$0.71 \times 0.47$	-31.00, -3.04*, -3.04*, -43.00
[18]	RT/Duroid 5880	1	31.73	2.75	$1.80 \times 0.63$	-31.00, -8.00, -10.00, -35.00
[19]	RT/Duroid 5870	0	10.71	3.47, 4.73 and 6.31	$0.12 \times 0.24$	-27.89, -4.16, -3.88, -11.20
<b>This work</b>	FR-4	2	35.21	2.4	$0.48 \times 0.60$	-25.00, -3.40, -3.47, -40.10

$\lambda_g$  is the guided wavelength of the substrate material.

The S-parameters are given at the lowest frequency for each reference, and the superscript \* denotes the results obtained from simulations.

Tab. 4. Comparison of key parameters of the proposed CMRC-based QWLT PD and selected PDs in previous research works.

## 4. Conclusions

In this paper, a novel PD is miniaturized using asymmetric CMRC structures with low-cost materials based on QWLTs. The results showcase a remarkable miniaturization of 35.21% compared with  $\lambda_g/4$  and approximately 50% compared to the length of the conventional PD at 2.4 GHz, while maintaining good performance. Comparisons between simulated and measured results demonstrate good agreement, with the isolation  $|S_{23}|$  and magnitude of reflection coefficient  $|S_{11}|$  below 20 dB. The insertion loss  $|S_{21}|$  and  $|S_{31}|$  are at 3.3 dB. In addition, the proposed PD model is simple and straightforward to fabricate. In the future, the research will explore other asymmetric CMRC structures for additional size reduction of various microwave devices.

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